

# 155(8): Relativistic Time Delay Calculations

This is just a trivial extension of the light deflection calculation, and the latter is accepted unconditionally, with the result that major errors are introduced in exactly the same way. The light deflection calculation is based on:

$$\frac{d\phi}{dr} = \frac{1}{r^2} \left( \frac{1}{b^2} - \left( 1 - \frac{r_0}{r} \right) \left( \frac{1}{a^2} + \frac{1}{r^2} \right) \right)^{-1/2} \quad (1)$$

where  $a = \frac{L}{mc}$ ,  $b = \frac{cL}{E}$ ,  $r_0 = \frac{2mGr}{c^2}$ .  $-(2)$

The time delay calculation is the trivial extension:

$$\frac{d\tau}{dr} = \frac{d\tau}{d\phi} \frac{d\phi}{dr} \quad (3)$$

$$L = m r^2 \frac{d\phi}{d\tau} \quad (4)$$

$$E = m c^2 \frac{dt}{d\tau} \left( 1 - \frac{r_0}{r} \right) \quad (5)$$

hence  $\frac{dt}{dr} = \frac{dt}{d\tau} \frac{d\tau}{dr} = \frac{dt}{d\tau} \frac{d\tau}{d\phi} \frac{d\phi}{dr} \quad (6)$

$$\frac{dt}{dr} = \frac{1}{cb} \left( 1 - \frac{r_0}{r} \right)^{-1} \left( \frac{1}{b^2} - \left( 1 - \frac{r_0}{r} \right) \left( \frac{1}{a^2} + \frac{1}{r^2} \right) \right)^{-1/2} \quad (7)$$

From eq. (1), the light deflection is:

$$\Delta\phi = 2 \int_{R_0}^{\infty} \frac{1}{r^2} \left( \frac{1}{b^2} - \left( 1 - \frac{r_0}{r} \right) \left( \frac{1}{a^2} + \frac{1}{r^2} \right) \right)^{-1/2} dr \quad (8)$$

and for eq. (7) the time delay is:

$$\Delta t = \frac{1}{c b} \int \left( 1 - \frac{r_0}{r} \right)^{-1} \left( \frac{1}{b^2} - \left( 1 - \frac{r_0}{r} \right) \left( \frac{1}{a^2} + \frac{1}{r^2} \right) \right)^{-1/2} dr \quad - (9)$$

The two integrals (8) and (9) are the same except that we use  $\frac{1}{r^2}$  and the other  $(1 - r/r_0)^{-1}$ .

As in UFT 150 and note 150(8), the assumption used by Eistein in eq. (8) was:

$$\frac{1}{a^2} = ? \quad 0 \quad - (10)$$

Eistein also assumed a circular orbit so:

$$\frac{1}{b^2} = ? \left( 1 - \frac{r_0}{r} \right) \left( \frac{1}{a^2} + \frac{1}{r^2} \right) \quad - (11)$$

in which case the denominator in eq. (8) is zero. In UFT 150 Eistein's own integral was evaluated by computer, giving a result that was different by a factor of a million.

Eq. (9) can be written as:

$$\Delta t = \frac{1}{c} \int \left( 1 - \frac{r_0}{r} \right)^{-1} \left( 1 - \left( 1 - \frac{r_0}{r} \right) \left( \frac{1}{a^2} + \frac{1}{r^2} \right) b^2 \right)^{-1/2} dr \quad - (12)$$

This is a very simple equation to evaluate by

computer.

The Shapiro time delay is saved directly in eq. (12), but again assumes zero planet mass:

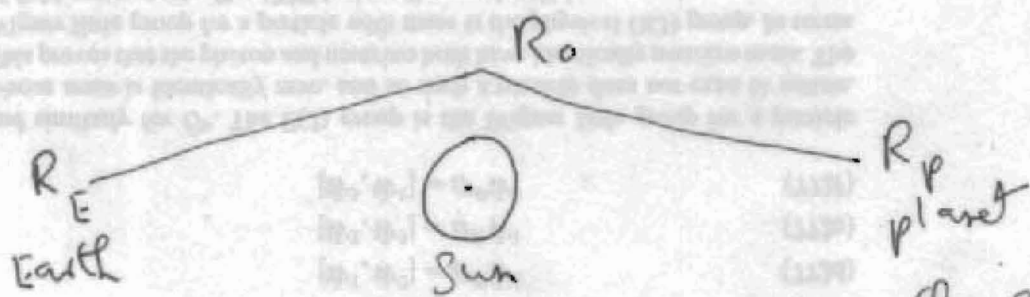
$$\frac{1}{a^2} = ? \quad 0 \quad - (13)$$

so Shapiro evaluated:

$$\Delta t = ? \quad \frac{1}{c} \int \left(1 - \frac{r_0}{r}\right)^{-1} \left(1 - \left(1 - \frac{r_0}{r}\right) \frac{b^2}{r^2}\right)^{-1/2} dr \quad - (14)$$

This is given in eq. (6.3.44) of R.M. Wald, "General Relativity" (Chicago, 1984). The notation used by Wald is:

$$\frac{r_0}{r} := \frac{2M}{r} \quad - (15)$$



A radar signal is emitted from the earth, located at the radial coordinate  $R_E$ . It grazes the sun at closest approach  $R_0$ , and is reflected off a planet at  $R_P$ .

The correct integral is therefore:



$$t = \frac{2}{c} \int_{R_E}^{R_0} \left(1 - \frac{r_0}{r}\right)^{-1} \left(1 - \left(1 - \frac{r_0}{r}\right) \left(\frac{1}{a^2} + \frac{1}{r^2}\right) b^2\right)^{-1/2} dr$$

$$+ \frac{2}{c} \int_{R_0}^{R_P} \left(1 - \frac{r_0}{r}\right)^{-1} \left(1 - \left(1 - \frac{r_0}{r}\right) \left(\frac{1}{a^2} + \frac{1}{r^2}\right) b^2\right)^{-1/2} dr \quad (16)$$

The factor 2 is due to the fact that the signal travels from  $R_E$  to  $R_P$  and then back from  $R_P$  to  $R_E$ .

Eq (16) is easily evaluated by computer.

The result claimed by Wald is:

$$\Delta t = \frac{2}{c} \left( (R_E^2 - R_0^2)^{1/2} + (R_P^2 - R_0^2)^{1/2} \right)$$

$$+ \frac{2GM}{c^3} \left( 2 \log_e \left( \frac{R_E + (R_E^2 - R_0^2)^{1/2}}{R_0} \right) \right. \quad (17)$$

$$\left. + 2 \log_e \left( \frac{R_P + (R_P^2 - R_0^2)^{1/2}}{R_0} \right) \right.$$

$$\left. + \left( \frac{R_E - R_0}{R_E + R_0} \right)^{1/2} + \left( \frac{R_P - R_0}{R_P + R_0} \right)^{1/2} \right)$$

from  $\frac{dt}{dr} = \left(1 - \frac{2m}{r}\right)^{-1} \left(1 - \left(1 - \frac{2m}{r}\right) \frac{b^2}{r^2}\right)^{-1/2} \quad (18)$

5) It is claimed by Wald that the result is obtained by integrating eq. (18) and then sic. "differentiating w.r. to  $M$  (holding  $R_0$  fixed) ... in close analogy to the derivation of the light bending effect."

It is claimed that null geodesics will be same  $R_0$  are obtained by varying  $M$  (the mass of the sun)

It is known from UFT 150 that the method

is wildly wrong.

Integral was evaluated, i.e.:

$$\Delta\phi = 2 \int_0^{1/R_0} (R_0^{-2} - 2MR_0^{-3} - u^2 + 2Mu^3)^{-1/2} du \quad (19)$$

The method used by Einstein is given in Wald, page 145, and is:

$$\begin{aligned} \left. \frac{d(\Delta\phi)}{dM} \right|_{M=0} &= 2 \left. \frac{\int_0^{1/R_0} (R_0^{-3} - u^3) du}{(R_0^{-2} - 2MR_0^{-3} - u^2 + 2Mu^3)^{3/2}} \right|_{M=0} \\ &= 2 \int_0^{1/b} \frac{(b^{-3} - u^3) du}{(b^{-2} - u^2)^{3/2}} \\ &= \frac{4}{b} \quad (20) \end{aligned}$$

The deflection of light to first order in  $M$

6)

is:

$$\delta\phi = \Delta\phi - \pi \sim M \left. \frac{d(\Delta\phi)}{dm} \right|_{m=0}$$

$$= \frac{4M}{b} - (21)$$

$$\delta\phi = \frac{4GM}{bc^2} - (22)$$

finally, it is assumed that:

$$b = R_0 - (23)$$

The computer evaluation of eq. (19) gives a totally different result, by six orders of magnitude.

The mass of the sun,  $M$ , is constant, it does not vary, so the method used in eq. (20) is meaningless. The same method is used to evaluate eq. (18), with the same assumption of zero photon mass:

$$\frac{1}{a^2} = ? \quad 0 - (24)$$

Therefore the result (17) is totally incorrect.

The correct result is obtained from evaluating eq. (16) by computer, with finite photon mass.