

172(10) : Derivation of the Thomas Factor and
 Spin Orbit Coupling for the Fermi Equation.

Start with the fermi equation in the format of
 eq. (10) of note 172(6):

$$((E - e\phi)^2 - m^2 c^4) \phi^R = c^2 (\underline{\sigma} \cdot \underline{\pi})(\underline{\sigma} \cdot \underline{\pi}) \phi^R \quad - (1)$$

where $\underline{\pi} = \underline{p} - e \underline{A} \quad - (2)$

Consider the case: $\underline{A} = \underline{0}, \phi \neq 0 \quad - (3)$

then:

$$\frac{1}{2m} (\underline{\sigma} \cdot \underline{p})(\underline{\sigma} \cdot \underline{p}) \phi^R = \frac{1}{2mc^2} ((E - e\phi)^2 - m^2 c^4) \phi^R$$

$$= \frac{1}{2mc^2} (c^2 p^2 - 2e\phi E + e^2 \phi^2) \phi^R \quad - (4)$$

$$E^2 - m^2 c^4 = p^2 c^2 \quad - (5)$$

using

$$\text{So: } \frac{p^2}{2m} \phi^R = \left(\frac{1}{2m} (\underline{\sigma} \cdot \underline{p})(\underline{\sigma} \cdot \underline{p}) + \frac{e\phi E}{mc^2} - \frac{e^2 \phi^2}{2mc^2} \right) \phi^R \quad - (6)$$

Now use $E = \gamma mc^2 \quad - (7)$

and consider the non-relativistic limit:
 $v \ll c, \quad - (8)$

in which $\gamma \rightarrow 1. \quad - (9)$

2) Then:

$$\frac{\mathbf{p}^2}{2m} \phi^R = \left(\frac{1}{2m} (\underline{\sigma} \cdot \underline{p})(\underline{\sigma} \cdot \underline{p}) + e\phi - \frac{e^2 \phi^2}{2mc^2} \right) \phi^R \quad - (10)$$

However:

$$\frac{\mathbf{p}^2}{2m} \phi^R = \frac{1}{2m} (\underline{\sigma} \cdot \underline{p})(\underline{\sigma} \cdot \underline{p}) \phi^R \quad - (11)$$

So

$$e\phi = 2mc^2 \quad - (12)$$

or

$$\frac{e\phi}{2mc^2} = 1 \quad - (13)$$

Therefore:

$$\frac{\mathbf{p}^2}{2m} \phi^R = \left(\frac{1}{2m} (\underline{\sigma} \cdot \underline{p}) \frac{e\phi}{2mc^2} (\underline{\sigma} \cdot \underline{p}) \right) \phi^R \quad - (14)$$

$$= \frac{e}{4m^2 c^2} \left((\underline{\sigma} \cdot \underline{p}) \phi (\underline{\sigma} \cdot \underline{p}) \right) \phi^R$$

Referring to notes 146(3), page 7, we have:

$$\begin{aligned} \underline{\sigma} \cdot \underline{p} \phi \underline{\sigma} \cdot \underline{p} &= \frac{\hbar}{i} \underline{\sigma} \cdot \underline{\nabla} (\phi \underline{\sigma} \cdot \underline{p}) \\ &= \frac{\hbar}{i} \underline{\sigma} \cdot \left((\underline{\nabla} \phi) \underline{\sigma} \cdot \underline{p} + \phi (\underline{\sigma} \cdot \underline{p})(\underline{\sigma} \cdot \underline{p}) \right) \\ &= -\frac{\hbar}{i} (\underline{\sigma} \cdot \underline{E})(\underline{\sigma} \cdot \underline{p}) + \phi \mathbf{p}^2 \end{aligned} \quad - (15)$$

3) where

$$\underline{E} = -\underline{\nabla} \phi \quad - (16)$$

is the electric field strength in volts m^{-1} .

Using the algebra of Pauli matrices:

$$(\underline{\sigma} \cdot \underline{E})(\underline{\sigma} \cdot \underline{p}) = \underline{E} \cdot \underline{p} + i \underline{\sigma} \cdot \underline{E} \times \underline{p} \quad - (17)$$

So:

$$\underline{\sigma} \cdot \underline{p} \phi \underline{\sigma} \cdot \underline{p} = i \hbar \underline{E} \cdot \underline{p} - \hbar \underline{\sigma} \cdot \underline{E} \times \underline{p} + \phi p^2 \quad - (18)$$

The spin-orbit term is:

$$\hat{H}_{so} = - \frac{e \hbar}{4 m^2 c^2} \underline{\sigma} \cdot \underline{E} \times \underline{p} \quad - (19)$$

Eq. (19) is usually written as:

$$\hat{H}_{so} = - \frac{e \underline{S} \cdot \underline{E} \times \underline{p}}{2 m^2 c^2} \quad - (20)$$

where

$$\underline{S} = \frac{1}{2} \hbar \underline{\sigma} \quad - (21)$$

The Thomas factor is the extra factor 2 in denominator of eq. (20).

4) If we consider an electric field of Q type:

$$\underline{E} = -\underline{\nabla} \phi = -\frac{e}{4\pi\epsilon_0} \frac{\underline{r}}{r^3}, \quad - (22)$$

$$\phi = -\frac{e}{4\pi\epsilon_0 r}, \quad - (23)$$

then: $\hat{H}_{so} = -\left(\frac{e^2}{4\pi\epsilon_0 r^3}\right) \frac{\underline{S} \cdot \underline{L}}{2m^2 c^2} \quad - (24)$

where $\underline{L} = \underline{r} \times \underline{p} \quad - (25)$

