

235(b): Spin Connection & Fundamental Rotation of Axes.

In the plane polar coordinate system:

$$\underline{e}_r = \underline{i} \cos \theta(t) + \underline{j} \sin \theta(t) \quad - (1)$$

$$\underline{e}_\theta = -\sin \theta(t) \underline{i} + \cos \theta(t) \underline{j} \quad - (2)$$

Therefore \underline{e}_r and \underline{e}_θ depend on time, and the unit vectors rotate. The spin connection defines this rotation.

In tensor notation the basis vectors are:

$$e^{(a)}_r = e^{(a)}, \quad e^{(a)}_\theta = e^{(a)} \quad - (3)$$

The covariant derivative is defined by:

$$D_\mu e^{(a)} = \partial_\mu e^{(a)} + \omega^{(a)}_{\mu(b)} e^{(b)} \quad - (4)$$

for example

$$\frac{D e^{(a)}}{dt} = \left(\frac{d e^{(a)}}{dt} \right)_{\text{static}} + \omega^{(a)}_{1(b)} e^{(b)} \quad - (5)$$

However:

$$\left(\frac{d e^{(a)}}{dt} \right)_{\text{static}} = 0 \quad - (6)$$

Because for static coordinates θ is not time dependent. So for θ coordinates:

$$\frac{D e^{(a)}}{dt} = \omega^{(a)}_{1(b)} e^{(b)} \quad - (7)$$

2) In vector notation:

$$\boxed{\frac{d\underline{e}_r}{dt} = \omega^{(1)}_{(2)} \underline{e}_\theta} \quad - (8)$$

This result is always denoted by:

$$\frac{d\underline{e}_r}{dt} = \omega \underline{e}_\theta \quad - (9)$$

but rigorously, it should be: $\quad - (10)$

$$\boxed{\frac{d\underline{e}_r(t)}{dt} = \left(\frac{d\underline{e}_r}{dt}\right)_{\text{static}} + \omega \underline{e}_\theta(t)}$$

It follows that:

$$\boxed{\omega^{(1)}_{(2)} = \omega} \quad - (11)$$

a result of basic importance. The fundamental equation (10) can be expressed as:

$$\frac{d\underline{e}_r(t)}{dt} = \frac{d\underline{e}_r(\text{static})}{dt} + \underline{\omega} \times \underline{e}_r \quad - (12)$$

if :

$$3) \quad \underline{\omega} = \omega \underline{k} \quad - (13)$$

and

$$\underline{k} = \underline{e}_r \times \underline{e}_\theta \quad - (14)$$

$$\underline{e}_\theta = \underline{k} \times \underline{e}_r \quad - (15)$$

The precession vector $\underline{\omega}$ is fundamental to all dynamics.

Similarly:

$$\frac{D\underline{r}}{dt} = \left(\frac{d\underline{r}}{dt} \right)_{\text{static}} + \underline{\omega} \times \underline{r} \quad - (16)$$

where

$$\underline{r} = r \underline{e}_r, \quad - (17)$$

$$\begin{aligned} \text{i.e. } \frac{D}{dt}(r \underline{e}_r) &= \left(\frac{d}{dt}(r \underline{e}_r) \right)_{\text{static}} + r \underline{\omega} \times \underline{e}_r \\ &= \frac{dr}{dt} \underline{e}_r + \underline{\omega} \times \underline{r} \quad - (18). \end{aligned}$$

In the theory of angular momentum \underline{L} :

$$\underline{L} = \underline{r} \times \underline{p} \quad - (19)$$

$$= m \underline{r} \times \underline{v}.$$

Here

$$\underline{v} = \frac{d\underline{r}}{dt} \quad - (20)$$

4) So Torque is therefore:

$$\underline{Tq} = \frac{d\underline{L}}{dt} = m \left(\frac{d\underline{r}}{dt} \times \underline{v} + \underline{r} \times \underline{a} \right) \quad -(21)$$

where

$$\underline{a} = \frac{d\underline{v}}{dt}$$

-(22)

$$\text{So } \underline{Tq} = \frac{d\underline{L}}{dt} = m \underline{r} \times \underline{a} \quad -(23)$$

The acceleration is:

$$\underline{a} = \frac{d^2 \underline{r}}{dt^2} \underline{e}_r + \frac{d\underline{\omega}}{dt} \times \underline{r} \quad -(24)$$

$$+ \underline{\omega} \times (\underline{\omega} \times \underline{r}) + 2 \underline{v}_s \times \underline{\omega}$$

$$= \left(\frac{d^2 r}{dt^2} - \omega^2 r \right) \underline{e}_r + \left(r \frac{d\omega}{dt} + 2\omega \frac{dr}{dt} \right) \underline{e}_\theta$$

In a static or vertical frame:

$$\underline{a} = \frac{d^2 r}{dt^2} \underline{e}_r \quad -(25)$$

So the torque is due entirely to:

$$\underline{Tq} = \frac{d\underline{L}}{dt} = m \underline{r} \times \left(r \frac{d\omega}{dt} + 2\omega \frac{dr}{dt} \right) \underline{e}_\theta$$

-(26)

5) i.e.

$$\underline{T}_Q = \frac{d\underline{L}}{dt} = m \left(r \frac{d\omega}{dt} + 2 \frac{dr}{dt} \omega \right) \underline{k} \quad - (27)$$
$$= 2(\underline{v}_s \times \underline{\omega}) \times \underline{r} + \underline{r} \times \left(\frac{d\omega}{dt} \times \underline{r} \right)$$

for motion in a plane.

For any orbit:

$$\underline{T}_Q = \underline{0} \quad - (28)$$

because \underline{J} is a constant of motion.

In general the angular momentum can be expressed as:

$$\underline{L} = m \underline{r} \times (\underline{\omega} \times \underline{r}) \quad - (29)$$
$$= m \left(r^2 \underline{\omega} - \underline{r} (\underline{r} \cdot \underline{\omega}) \right)$$

If: $\underline{\omega} = \omega \underline{k}, \underline{r} = r \underline{e}_r \quad - (30)$

then

$$\underline{L} = m r^2 \underline{\omega} \quad - (31)$$

for motion in a plane. Observe $\underline{L} \neq \underline{\omega}$.

For motion in a plane the angular momentum is the
specification vector.