

## 240(1): Development of the Verrier Calculation by Fawcett

This is based on the gravitational potential:

$$\phi(\underline{r}) = -G \sum_{i=1}^N \frac{m_i}{|\underline{r}_i - \underline{r}|} \quad (1)$$

which is purely Newtonian. However, if GR is to be applied correctly it should be based on:

$$\phi(\underline{r}) = -G \sum_{i=1}^N \left( \frac{m_i}{|\underline{r}_i - \underline{r}|} \left( 1 + \frac{dm_i G}{c^2 |\underline{r}_i - \underline{r}|} \right) \right) \quad (2)$$

as in note 239(a) and UFT 239.

Here:

$$\frac{1}{|\underline{r}' - \underline{r}|} = (r^2 - 2 \underline{r} \cdot \underline{r}' + r'^2)^{-1/2} \quad (3)$$

For a continuous distribution:

$$\phi(\underline{r}) = -G \int \frac{\rho(\underline{r}')}{|\underline{r}' - \underline{r}|} d^3 r' \quad (4)$$

for eq. (1), but this should be:

$$\phi(\underline{r}) = -G \int \frac{\rho(\underline{r}')}{|\underline{r}' - \underline{r}|} \left( 1 + \frac{d\rho(\underline{r}') G}{c^2 |\underline{r}' - \underline{r}|} \right) d^3 r' \quad (5)$$

For self consistency, eq. (5) must be used in

2) The gravitational N body problem.

The procedure used by Farside is to calculate:

$$\phi(r, \theta) = -G \int_0^\infty \int_0^\pi \int_0^{2\pi} \frac{r'^2 \rho(r', \theta') \sin \theta'}{|r - r'|} dr' d\theta' d\phi' \quad (6)$$

$$= -2\pi G \int_0^\infty \int_0^\pi r'^2 \rho(r', \theta') \sin \theta' \left\langle \frac{1}{|r - r'|} \right\rangle dr' d\theta'$$

where:

$$\left\langle \frac{1}{|r - r'|} \right\rangle = \frac{1}{r} \sum_{n=0}^{\infty} \left( \frac{r'}{r} \right)^n P_n(\cos \theta) P_n(\cos \theta') \quad (7)$$

when

$$r > r' \quad (8)$$

and  $\left\langle \frac{1}{|r - r'|} \right\rangle = \frac{1}{r'} \sum_{n=0}^{\infty} \left( \frac{r}{r'} \right)^n P_n(\cos \theta) P_n(\cos \theta')$

when

$$r < r' \quad (9)$$

These are well known methods involving Legendre polynomials  $P_n(\cos \theta)$ . The volume integral is:

$$\int d^3 r' = \int_0^\infty \int_0^\pi \int_0^{2\pi} r'^2 \sin \theta' dr' d\theta' d\phi' \quad (10)$$

Eq. (6) is purely Newtonian, but if EGR is taken seriously it should be as follows:



$$\phi(r, \theta) = -G \int_0^\infty \int_0^\pi \int_0^{2\pi} r'^2 \left( \frac{\rho(r')}{|\underline{r}' - \underline{r}|} \left( 1 + \frac{d\rho(r')}{dr'} G \right) \right) \sin\theta' dr' d\theta' d\phi' \quad - (11)$$

and this introduces profound differences into the whole theory.

The Farside theory considers a uniform ring of mass  $M$  and radius  $r$  centred at the origin, but completely neglects the second term in eq. (11). The Farside theory produces:

$$\phi(r) = -\frac{MG}{r} \left( 1 + \frac{1}{4} \left( \frac{a}{r} \right)^2 + \frac{9}{64} \left( \frac{a}{r} \right)^4 + \dots \right) \quad \text{for } r > a \quad - (12)$$

$$\text{and } \phi(r) = -\frac{MG}{a} \left( 1 + \frac{1}{4} \left( \frac{r}{a} \right)^2 + \frac{9}{64} \left( \frac{r}{a} \right)^4 + \dots \right) \quad \text{for } r < a \quad - (13)$$

$$\text{for } r < a \quad - (14)$$

These results are applied to the solar system to produce the result for force per unit mass on a planet due to the sun and all the other planets. The mass of the sun is denoted  $M_0$ . With these definitions the result is:

$$4) F(R_i) = -\frac{M_0 G}{R_i^2} - \frac{G}{R_i^2} \sum_{j < i} m_j \left( 1 + \frac{3}{4} \left( \frac{R_i}{R_j} \right)^2 + \frac{45}{64} \left( \frac{R_i}{R_j} \right)^4 + \dots \right) \\ + \frac{G}{R_i^2} \sum_{j > i} m_j \left( \frac{R_i}{R_j} \right) \left( \frac{1}{2} \left( \frac{R_i}{R_j} \right)^2 + \frac{9}{16} \left( \frac{R_i}{R_j} \right)^4 + \dots \right) \quad - (16)$$

In order to simplify the problem consider:  
 $j < i, \quad j = 1, \quad i = 2 \quad - (17)$

Then:

$$F(R_2) = -\frac{M_0 G}{R_2^2} - \frac{G}{R_2^2} m_1 \left( 1 + \frac{3}{4} \left( \frac{R_1}{R_2} \right)^2 + \dots \right) \quad - (18)$$

$$= -\frac{(M_0 + m_1)G}{R_2^2} - \frac{3}{4} \frac{m_1 G R_1^2}{R_2^4} \quad - (19)$$

In a rough first approximation add the term due to the EGR theory for the sun, then:

$$F(R_2) = -\frac{(M_0 + m_1)G}{R_2^2} - \frac{3}{2} \frac{G}{R_2^4} \left( \frac{m_1 R_1^2}{2} + \alpha M_0 r_0 \right) \quad - (20)$$

It is seen that the term coming from



Precession is changed as follows:

$$\frac{1}{2} M_1 R_1^2 \rightarrow \frac{1}{2} M_1 R_1^2 + 2 M_0 r_0 \quad - (21)$$

If:

$$R_1 = \text{sun to mercury distance} = 5.79 \times 10^{10} \text{ m}$$

$$R_2 = \text{sun to earth distance} = 1.50 \times 10^{11} \text{ m}$$

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$$M_0 = \text{mass of sun} = 2 \times 10^{30} \text{ kg}$$

$$r_0 = \frac{2 M_0 G}{c^2} = 2.950 \times 10^4 \text{ m}$$

$$d(\text{Earth}) = 1.496 \times 10^{11} \text{ m}$$

$$M_1 = \text{mass of Mercury} = 3.3 \times 10^{23} \text{ kg}$$

$$\text{Then: } \frac{1}{2} M_1 R_1^2 = \frac{3.3}{2} \times 10^{23} \times 5.79^2 \times 10^{20}$$

$$= 5.53 \times 10^{44} \text{ kg m}^2$$

$$\text{and } 2 M_0 r_0 = 1.496 \times 10^{11} \times 2 \times 10^{30} \times 2.95 \times 10^4$$

$$= 8.83 \times 10^{45} \text{ kg m}^2$$

$$\text{The precession term is almost doubled by the sun's effect alone. Einstein completely rejected this effect and his theory cannot be true.}$$