

353(5): Development of Wave Equations

Consider the definition of \underline{E}_F and \underline{B}_F in the generalized Kramé theory:

$$\underline{E}_F = -\underline{\nabla} \Phi - \frac{\partial \underline{v}}{\partial t} \quad - (1)$$

$$\underline{B}_F = \underline{W} = \underline{\nabla} \times \underline{v} \quad - (2)$$

where:

$$\underline{\nabla} \Phi = \underline{\nabla} (h + \phi) - \underline{f}_{\text{visc}} \quad - (3)$$

and where the viscous force is defined by:

$$\underline{f}_{\text{visc}} = \mu \nabla^2 \underline{v} + (\mu + \mu') \underline{\nabla} (\underline{\nabla} \cdot \underline{v}) \quad - (4)$$

The generalized Kramé equations are then equivalent to the rigorous Navier Stokes equation and are:

$$\underline{\nabla} \cdot \underline{E}_F = \rho_F \quad - (5)$$

$$\underline{\nabla} \times \underline{E}_F + \frac{\partial \underline{B}_F}{\partial t} = \underline{0} \quad - (6)$$

$$\underline{\nabla} \cdot \underline{B}_F = 0 \quad - (7)$$

and

$$\underline{\nabla} \times \underline{B}_F - \frac{1}{a_0^2} \frac{\partial \underline{E}_F}{\partial t} = \frac{1}{a_0^2} \underline{J}_F \quad - (8)$$

where

$$\rho_F = \underline{\nabla} \cdot ((\underline{v} \cdot \underline{\nabla}) \underline{v}) \quad - (9)$$

and

$$\underline{J}_F = a_0^2 \underline{\nabla} \times (\underline{\nabla} \times \underline{v}) - \frac{\partial}{\partial t} ((\underline{v} \cdot \underline{\nabla}) \underline{v}) \quad - (10)$$

Eqs. (6) and (7) follow directly from eqs.

2) (1) and (2):

$$\underline{\nabla} \cdot \underline{B}_F = \underline{\nabla} \cdot \underline{\nabla} \times \underline{V} \equiv 0 \quad - (11)$$

and

$$\begin{aligned} \underline{\nabla} \times \left(-\underline{\nabla} \Phi - \frac{\partial \underline{V}}{\partial t} \right) + \frac{\partial}{\partial t} (\underline{\nabla} \times \underline{V}) & - (12) \\ = -\underline{\nabla} \times \underline{\nabla} \Phi - \frac{\partial}{\partial t} (\underline{\nabla} \times \underline{V}) + \frac{\partial}{\partial t} \underline{\nabla} \times \underline{V} & = 0 \end{aligned}$$

Eq. (5) is:

$$\underline{\nabla} \cdot \left(-\underline{\nabla} \Phi - \frac{\partial \underline{V}}{\partial t} \right) = \rho_F \quad - (13)$$

i.e.

$$\boxed{\nabla^2 \Phi + \frac{\partial}{\partial t} (\underline{\nabla} \cdot \underline{V}) = -\rho_F} \quad - (14)$$

This should be compared with Eq. (8) of Note 352(6):

$$\nabla^2 \phi_w + \frac{\partial}{\partial t} (\underline{\nabla} \cdot \underline{W}) = -\rho / \epsilon_0 \quad - (15)$$

Eq. (8) becomes:

$$\begin{aligned} \underline{\nabla} \times (\underline{\nabla} \times \underline{V}) - \frac{1}{a_0^2} \frac{\partial}{\partial t} \left(-\frac{\partial \underline{V}}{\partial t} - \underline{\nabla} \Phi \right) & = \frac{\underline{J}_F}{a_0^2} \\ = -\nabla^2 \underline{V} + \underline{\nabla} (\underline{\nabla} \cdot \underline{V}) + \frac{1}{a_0^2} \frac{\partial^2 \underline{V}}{\partial t^2} + \frac{1}{a_0^2} \underline{\nabla} \Phi & - (16) \end{aligned}$$

so

$$\boxed{\square \underline{V} + \underline{\nabla} \left(\underline{\nabla} \cdot \underline{V} + \frac{1}{a_0^2} \frac{\partial \Phi}{\partial t} \right) = \frac{1}{a_0^2} \underline{J}_F} \quad - (17)$$

here

$$\square := \frac{1}{a_0^2} \frac{\partial^2}{\partial t^2} - \nabla^2 \quad - (18)$$

3) Define

$$V^\mu := \left(\frac{\Phi}{a_0}, \underline{V} \right) \quad - (19)$$

and

$$d_\mu := \left(\frac{1}{a_0} \frac{\partial}{\partial t}, \underline{\nabla} \right) \quad - (20)$$

It follows that:

$$d_\mu V^\mu = \frac{1}{a_0^2} \frac{\partial \Phi}{\partial t} + \underline{\nabla} \cdot \underline{V} \quad - (21)$$

If the Lorenz gauge is assumed:

$$d_\mu V^\mu = 0 \quad - (22)$$

then Eq. (17) reduces to:

$$\square \underline{V} = \frac{1}{a_0^2} \underline{J}_F \quad - (23)$$

Eqs. (17) and (23) have the structure of the ECE wave equation.

If the Lorenz gauge (22) is used in Eq (14), it follows that:

$$\boxed{\square \Phi = \nabla_F} \quad - (24)$$

i.e.

$$\left(\frac{1}{a_0^2} \frac{\partial^2}{\partial t^2} - \nabla^2 \right) \Phi = \nabla_F \quad - (25)$$

Now define:

$$T^{\mu} = (a_0 q_F, \underline{T}_F) - (26)$$

then

$$\boxed{\square \nabla^{\mu} = T^{\mu}} - (27)$$

which can be regarded as the second order wave equation of fluid mechanics.

Eqs. (9) and (10) imply the continuity equation:

$$\frac{\partial q_F}{\partial t} + \underline{\nabla} \cdot \underline{T}_F = 0 - (28)$$

i.e

$$\partial_{\mu} T^{\mu} = 0 - (29)$$

Eq. (29) gives a justification for the Lorenz gauge (22) if:

$$T^{\mu} \propto \nabla^{\mu} - (30)$$

i.e

$$\underline{T}_F \propto \underline{\nabla} - (31)$$

and

$$q_F \propto \Phi - (32)$$

The units of Φ are $m^2 s^{-2}$ and those of q_F are s^{-2} . The units of $\underline{\nabla}$ are $m s^{-1}$ and of units of \underline{T}_F are $m s^{-3}$, so it follows that one possibility

is:

$$\nabla^2 \Phi = q_F - (33)$$

and

$$\frac{\partial^2 \underline{\nabla}}{1 \text{ m}^2} = \underline{T}_F - (34)$$

However, another possibility is:

which is eq. (25), and:

$$\square \Phi = \nabla F - (35)$$

$$\square \underline{v} = \frac{\underline{J}}{a_0^2} - (36)$$

which is eq. (23).

Therefore eqs. (22), (24) and (31) are consistent with the continuity equation and Lorenz gauge. QED.

Units Check

$$\square \Phi = m^{-2} m^2 s^{-2} = s^{-2}, \nabla F = s^{-2} \checkmark$$

$$\square \underline{v} = m^{-2} m s^{-1} = m^{-1} s^{-1};$$

$$\frac{1}{a_0^2} \underline{J} = s^2 m^{-2} m s^{-3} = m^{-1} s^{-1} \checkmark$$

Conclusion: The wave equations are:

$$\square \Phi = \nabla F - (37)$$

$$\square \underline{v} = \frac{1}{a_0^2} \underline{J} - (38)$$

These are rigorously compatible with Navier-Stokes and continuity equations.