

# 425(4) : Basic Theory and Orbital Hamiltonian Equations

As Maria and Theresa chapter 6, the Hamiltonian is defined by the Euler Lagrange equations using:

$$p = \frac{\partial L}{\partial \dot{q}}, \quad \dot{p} = \frac{\partial L}{\partial q} \quad - (1)$$

Using:

$$\frac{\partial L}{\partial t} = 0 \quad - (2)$$

then

$$\begin{aligned} \frac{dL}{dt} &= \frac{\partial L}{\partial q} \dot{q} + \frac{\partial L}{\partial \dot{q}} \ddot{q} \quad - (3) \\ &= \dot{q} \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}} \right) + \frac{\partial L}{\partial \dot{q}} \ddot{q} \end{aligned}$$

Using:

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}} \right) = \frac{\partial L}{\partial q} \quad - (4)$$

it follows that:

$$\frac{d}{dt} \left( L - \dot{q} \frac{\partial L}{\partial \dot{q}} \right) = 0 \quad - (5)$$

The Hamiltonian is defined as:

$$H = \dot{q} \frac{\partial L}{\partial \dot{q}} - L \quad - (6)$$

so

$$\frac{dH}{dt} = 0 \quad - (7)$$

Therefore:

$$H(q, p, t) = p \dot{q} - L(q, \dot{q}, t) \quad - (8)$$

The Hamiltonian is a function of  $q$  and  $p$ , which are canonically conjugate variables:

$$\frac{\partial p}{\partial q} = 0 \quad - (9)$$

and the Lagrangian is a function of  $q$  and  $\dot{q}$ , which  $\dot{q}$  is the time derivative of  $q$ , so  $q$  and  $\dot{q}$  are not independent:

$$\frac{d\dot{q}}{dq} \neq 0. \quad (8)$$

The Hamilton equations are derived by considering:

$$dH = \frac{\partial H}{\partial q} dq + \frac{\partial H}{\partial p} dp + \frac{\partial H}{\partial t} dt \quad (9)$$

and

$$H = p\dot{q} - L \quad (10)$$

so

$$dH = \dot{q} dp + p d\dot{q} - \frac{\partial L}{\partial q} dq - \frac{\partial L}{\partial \dot{q}} d\dot{q} \quad (11)$$

From eqs (9) and (11):

$$\dot{q} = \frac{\partial H}{\partial p}, \quad \dot{p} = -\frac{\partial H}{\partial q} \quad (12)$$

which are the Hamilton equations, Q.E.D.

These are powerful first order differential equations used in many fields of physics. They are written in terms of generalized coordinates  $p$  and  $q$ .

The canonically conjugate variables  $p$  and  $q$  must be defined before the Hamilton equations can be applied. This can be done using the Lagrange equations or by inspection.

For the classical gravitational problem:

$$L = \frac{p^2}{2m} + \frac{-GMm}{r} \quad (13)$$

and in plane polar coordinates  $(r, \phi)$ :

$$L = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\phi}^2) + \frac{nmg}{r} \quad - (14)$$

with:

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{r}} = \frac{\partial L}{\partial r} \quad - (15)$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\phi}} = \frac{\partial L}{\partial \phi} \quad - (16)$$

The generalized momenta are:

$$p_r = p = \frac{\partial L}{\partial \dot{r}} \quad - (17)$$

$$p_\phi = L = \frac{\partial L}{\partial \dot{\phi}} \quad - (18)$$

and the generalized  $q$  are:

$$q_r = r, \quad q_\phi = \phi \quad - (19)$$

The Hamiltonian is:

$$H = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\phi}^2) - \frac{nmg}{r} \quad - (20)$$

From eq. (17):

$$p_r = m \dot{r}, \quad - (21)$$

and from eq. (18):

$$p_\phi = L = n r^2 \dot{\phi} \quad - (22)$$

s.

$$H = \frac{1}{2} m (p_r \dot{r} + p_\phi \dot{\phi}) - \frac{nmg}{q_r} \quad - (23)$$

$$= \frac{p_r^2}{2m} + \frac{p_\phi^2}{2mr^2} - \frac{nmg}{q_r}$$

The first Hamilton equation:

$$\dot{q}_r = \frac{\partial H}{\partial p_r} \quad - (24)$$

7) gives

$$\dot{q}_r = \frac{\partial H}{\partial p_r} = \frac{p_r}{m} \quad - (25)$$

i.e

$$\boxed{p_r = m \dot{r}} \quad - (26)$$

The first Hamilton equation also gives:

$$\dot{\phi} = \dot{q}_\phi = \frac{\partial H}{\partial p_\phi} = \frac{p_\phi}{mr^2} \quad - (27)$$

So:

$$\boxed{p_\phi = mr^2 \dot{\phi}} \quad - (28)$$

Eq.s. (26) and (28) mean:

$$\boxed{p = m \dot{r}, \quad L = mr^2 \dot{\phi}} \quad - (29)$$

The second Hamilton equation:

$$\dot{p} = - \frac{\partial H}{\partial q} \quad - (30)$$

gives:

$$\boxed{\dot{p}_r = - \frac{nmG}{r^2}} \quad - (31)$$

and

$$\boxed{\dot{p}_\phi = \dot{L} = 0} \quad - (32)$$

Eq. (31) means:

$$m \ddot{r} = - \frac{nmG}{r^2} \quad - (33)$$

which is the orbital force equation in vertical coordinates.

5) The vector function of  $\underline{e}_r$  (33) is:

$$m \ddot{\underline{r}} = -\frac{nm\hbar^2}{r^2} \underline{e}_r \quad - (34)$$

where  $\underline{e}_r$  is the radial unit vector. In plane polar coordinates:

$$\ddot{\underline{r}} = (\ddot{r} - r\dot{\phi}^2) \underline{e}_r + (r\ddot{\phi} + 2\dot{r}\dot{\phi}) \underline{e}_\phi \quad - (35)$$

so eq. (34) splits into:

$$m(\ddot{r} - r\dot{\phi}^2) = -\frac{nm\hbar^2}{r^2} \quad - (36)$$

and

$$r\ddot{\phi} + 2\dot{r}\dot{\phi} = 0 \quad - (37)$$

Eq. (37) is equivalent to:

$$\frac{dL}{dt} = 0 \quad - (38)$$

and eq. (36) can be derived from:

$$\frac{dH}{dt} = 0 \quad - (39)$$

Eq. (38) and (39) are from Eckardt equation

Note that

$$\dot{p} = \frac{\partial \mathcal{L}}{\partial r} = -\frac{\partial H}{\partial r} \quad - (40)$$

$$\frac{\partial(p\dot{r})}{\partial r} = 0 \quad - (41)$$

so

which implies that

$$\frac{\partial \dot{r}}{\partial r} = 0 \quad - (42)$$

This analysis must now be repeated for the relativistic and many